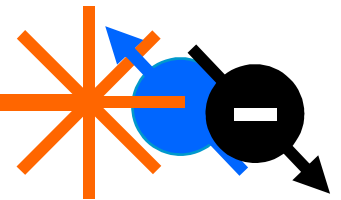


Ultrafast Coherent Electron Spin Flip in a 2D Electron Gas



Carey Phelps¹, Timothy Sweeney¹, Ronald T. Cox², Hailin Wang¹

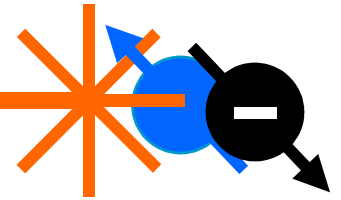
¹Department of Physics, University of Oregon, Eugene, OR 97403

²Nanophysics and Semiconductor Team, the Néel Institute and the CEA, Grenoble, France

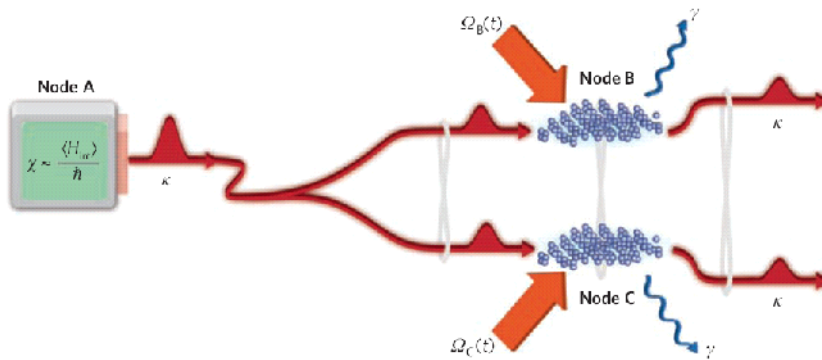


Supported by NSF and ARL

Electron Spins in Semiconductors



Ensemble spin Systems



Quantum memories and repeaters,
spintronic devices

Duan *et al.*, *Nature* **414** (2001).

Briegel *et al.*, *PRL* **81** (1998).

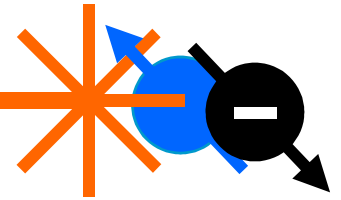
For recent reviews, see for example,
Kimble, *Nature* **453**, 1023 (2008).

Two goals:

- Complete spin control.
- Protect electron spins from decoherence.

→ Key step: π -rotations

Dynamic Decoupling



- **Uhrig dynamic decoupling (UDD)**

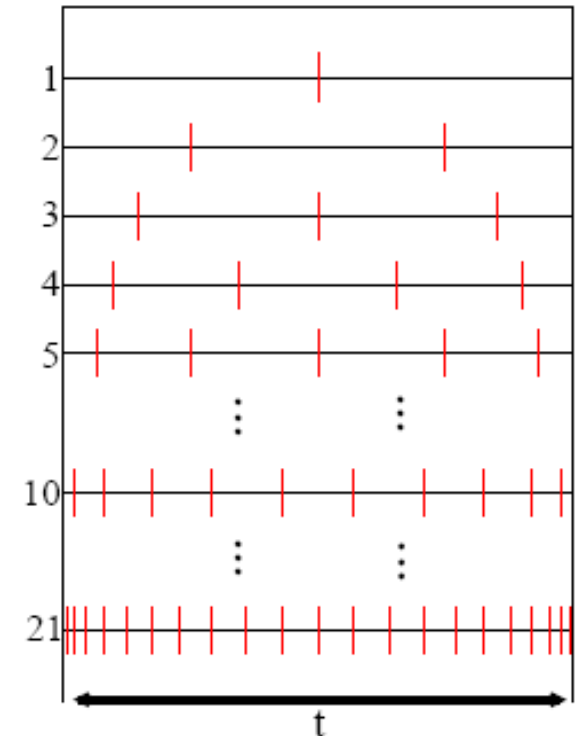
- Special sequence of π -pulses.
- Reverses decoherence due to dynamical interactions with the environment.

$$t_j = \frac{T}{2} \left(1 - \cos \frac{\pi j}{n+1}\right), \quad j=1, 2, \dots, n$$

n : number of pulses

- **UDD is “universal” and “magical.”**

- **Universal:** *Independent* of dephasing Hamiltonian.
- **Magical:** The improvement in decoherence time scales *linearly*, instead of exponentially, with the number of pulses.



G. S. Uhrig, *Phys. Rev. Lett.* **98** (2007).

B. Lee, W.M. Witzel, S. Das Sarma, *Phys. Rev. Lett.* 100, 160505 (2008).

W. Yao, R.B. Liu, and L. J. Sham, *Phys. Rev. Lett.* **98**, 077602 (2007).

Dynamic Decoupling

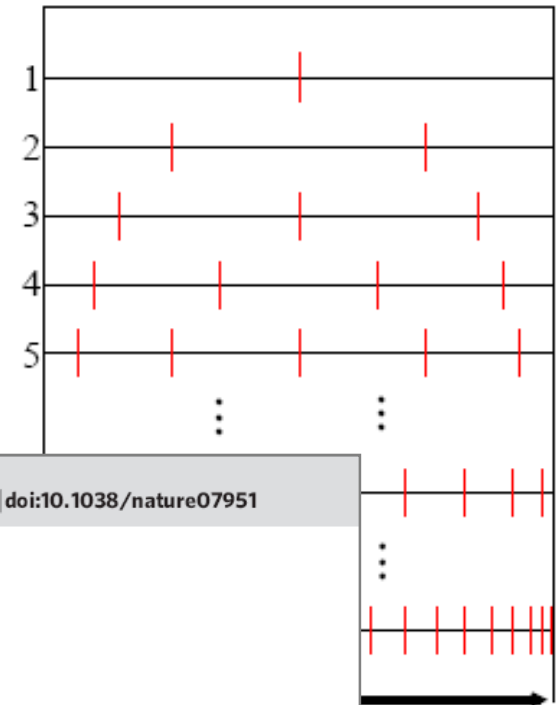


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nature

Vol 458 | 23 April 2009 | doi:10.1038/nature07951

LETTERS

Optimized dynamical decoupling in a model quantum memory

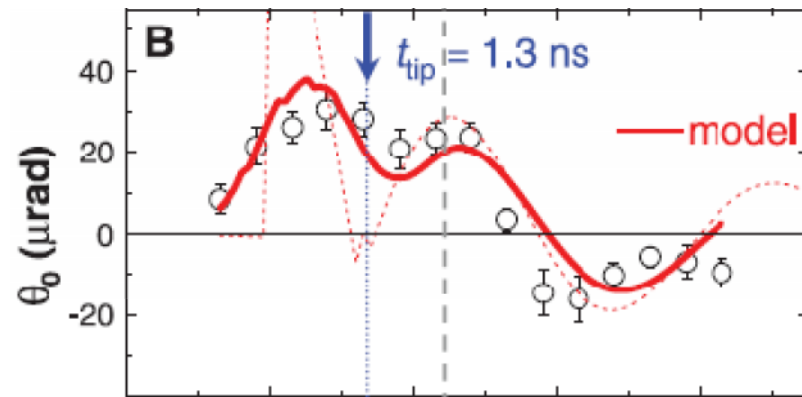
Michael J. Biercuk^{1,2*}, Hermann Uys^{1,3*}, Aaron P. VanDevender¹, Nobuyasu Shiga^{1†}, Wayne M. Itano¹ & John J. Bollinger¹

Single Quantum Dots

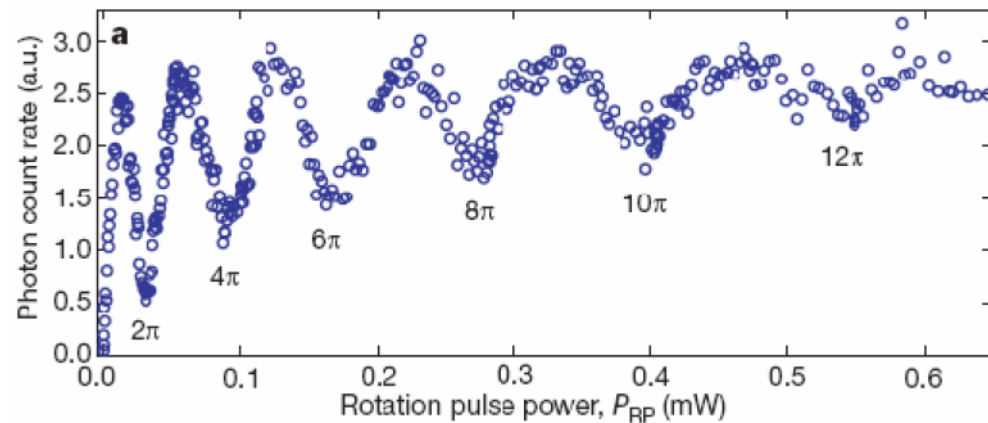


Complete spin control has been realized recently in single quantum dots:

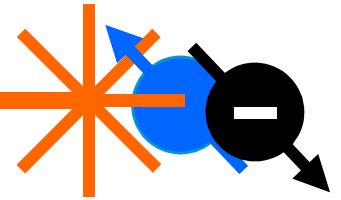
J. Berezovsky *et al.*, *Science* **320**, 349 (2008)



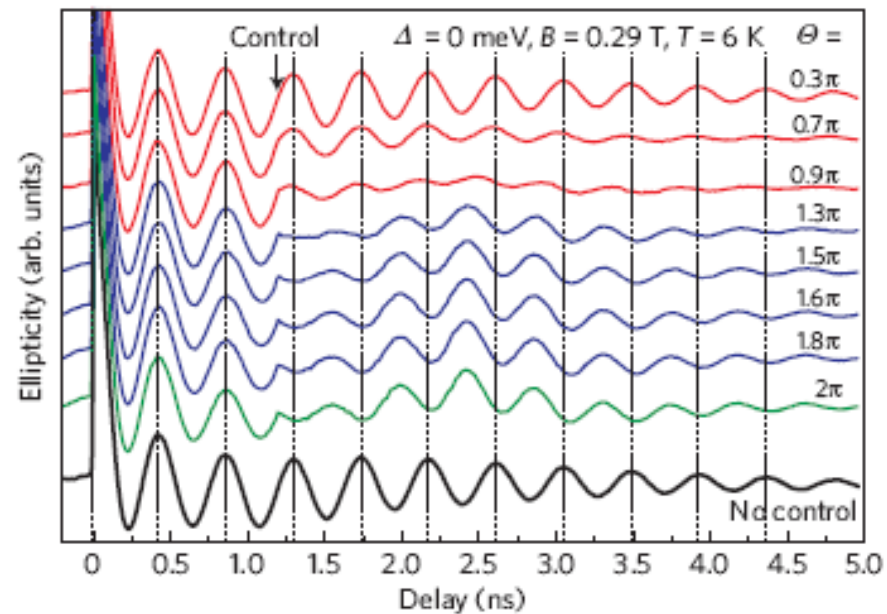
D. Press, T. D. Ladd, B. Zhang, and Y. Yamamoto, *Nature* **456** 218 (2008)



Quantum Dot Ensembles

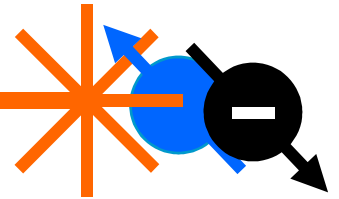


A. Greilich *et al.*, *Nature Physics* **5**, 262 (2009)



Quantum dots feature atomic-like, sharp optical transitions. Complete spin control in these systems is easier to achieve than in a 2DEG.

Spin Ensembles: 2D Electron Gas



- **Difficulties and physics issues** for π -rotation in 2D electron gas:

Inherent manybody interactions.

Short dipole decoherence time.

Quantum dots ~ 100 ps.

2D electron gas ~ 2 ps.

Earlier studies in 2DEGs have only achieved partial spin rotation.

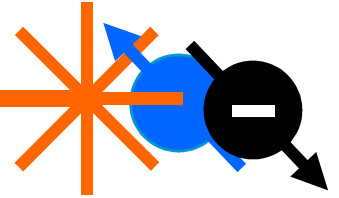
J. Gupta *et al.*, *Science* **292** (2001).

Y. Shen *et al.*, *Phys. Rev. B* **75** (2007).

Y. Wu *et al.*, *Phys. Rev. Lett* **99** (2007).

K. Fu *et al.*, *Nature Phys.* **4** (2008).

Outline

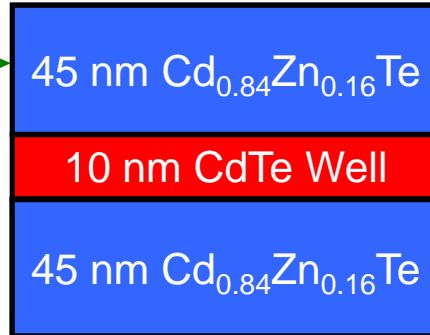


- Mechanism for spin rotation
- Experimental design
- Experimental results
- Summary

2D Electron Gas and Trions

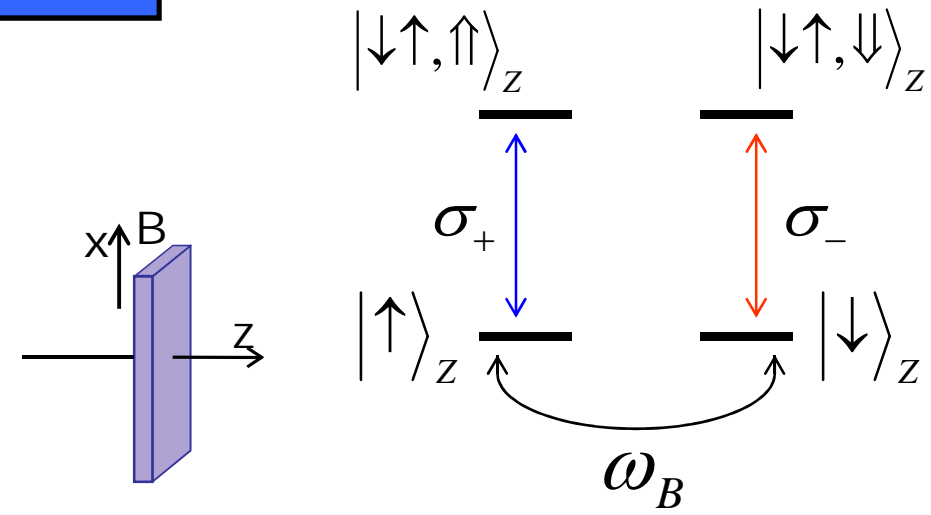
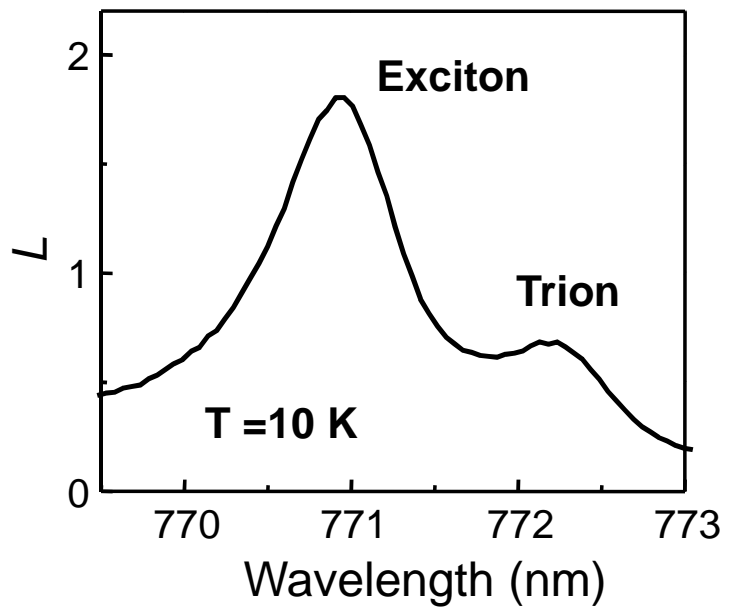


Indium doping
 $n \sim 3 \times 10^{10} / \text{cm}^2$



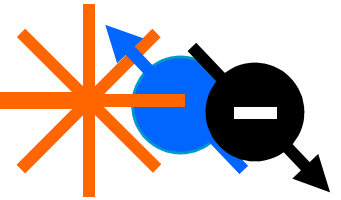
Sample from R. Cox,
 Grenoble, France

- Trions consist of two electrons of opposite spins bound to a hole.
- Trion absorption scales with electron doping density.

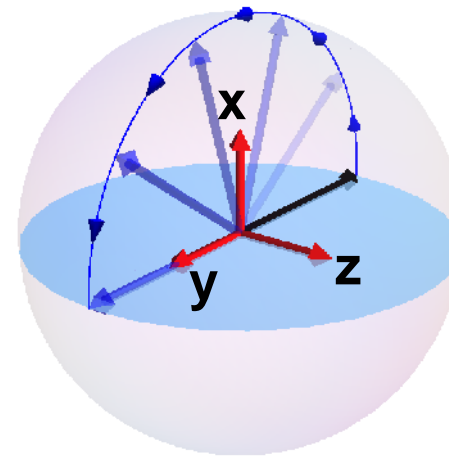
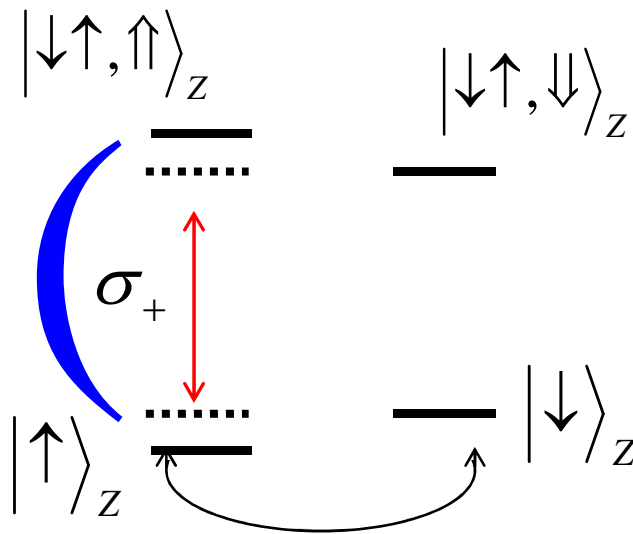


Magnetic field along x-direction leads to spin precession in y,z-plane (.45 T).

Spin Rotation via Optical Stark Effect



Pulse duration \ll Larmor precession period.



$$\phi = \int \delta\omega(t) dt$$

"Control" beam:

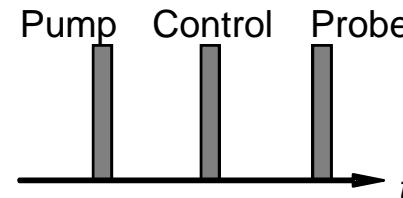
Induces a Stark shift, resulting in a rotation angle about the z-axis.

J.A. Gupta *et al.*, *Science* **292** (2001)

Experimental Setup

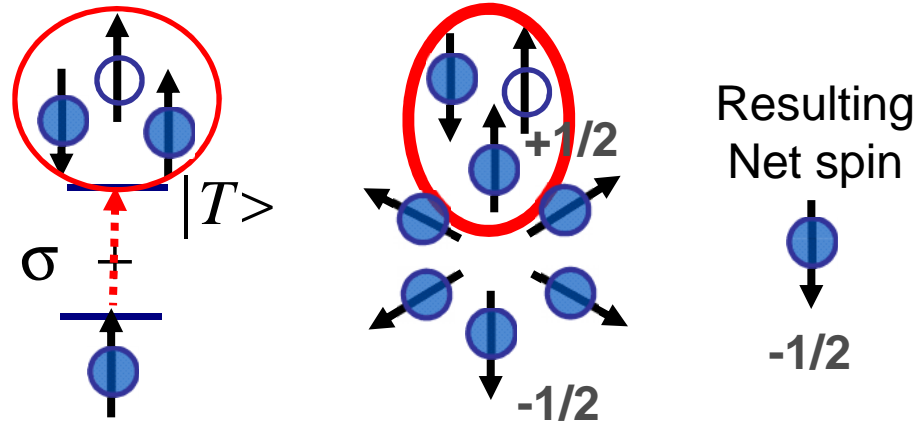


3 beam experiment



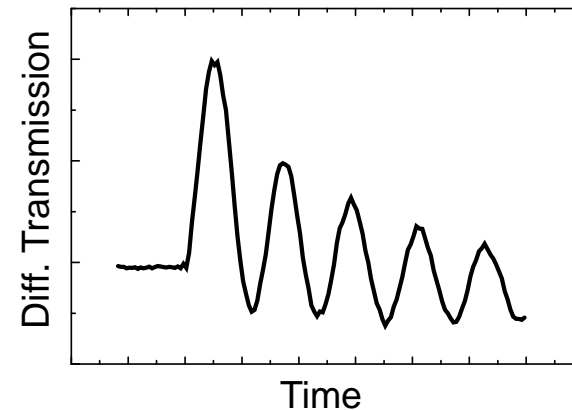
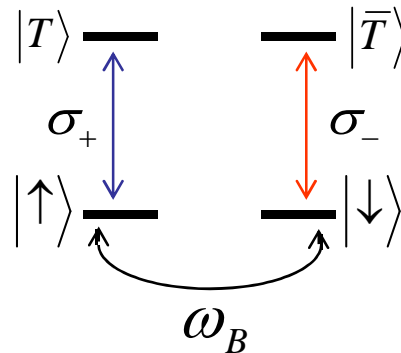
Spin initialization "Pump"

Trion excitation leads to a net spin polarization in the electron spin bath.

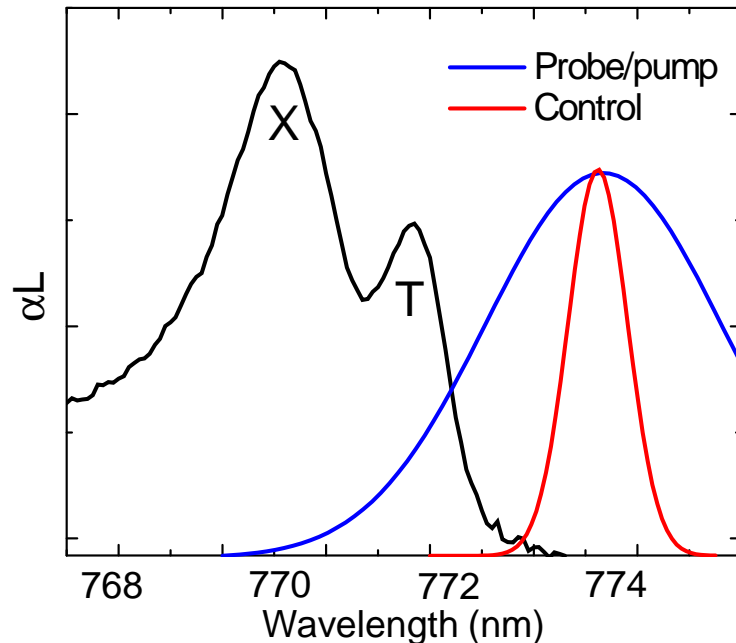
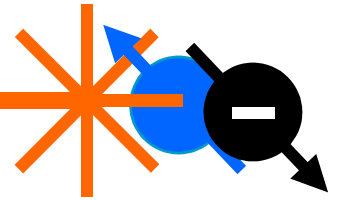


Measurement "Probe"

Transmission follows the resulting Larmor precession.



Experimental Design



3 control pulse parameters:

Δ , τ , and Ω .

τ no more than ~ 2 ps.

Ω less than trion binding energy.

These requirements conflict with (3).
Compromise: choose $\tau = 2$ ps and Δ small.

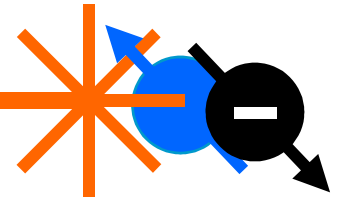
3 Goals

- 1.) Avoid dipole decoherence
- 2.) Avoid exciton excitation (many-body effects)
- 3.) $\phi = \int \delta\omega(t) dt \geq \pi$

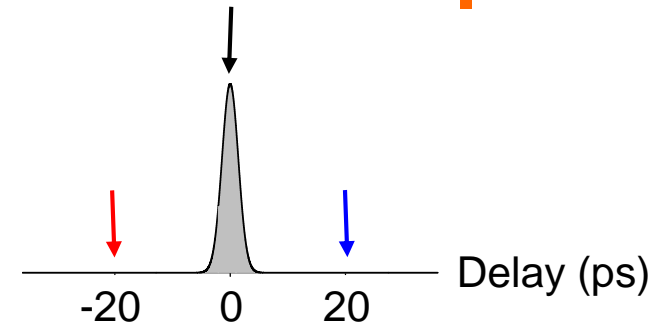
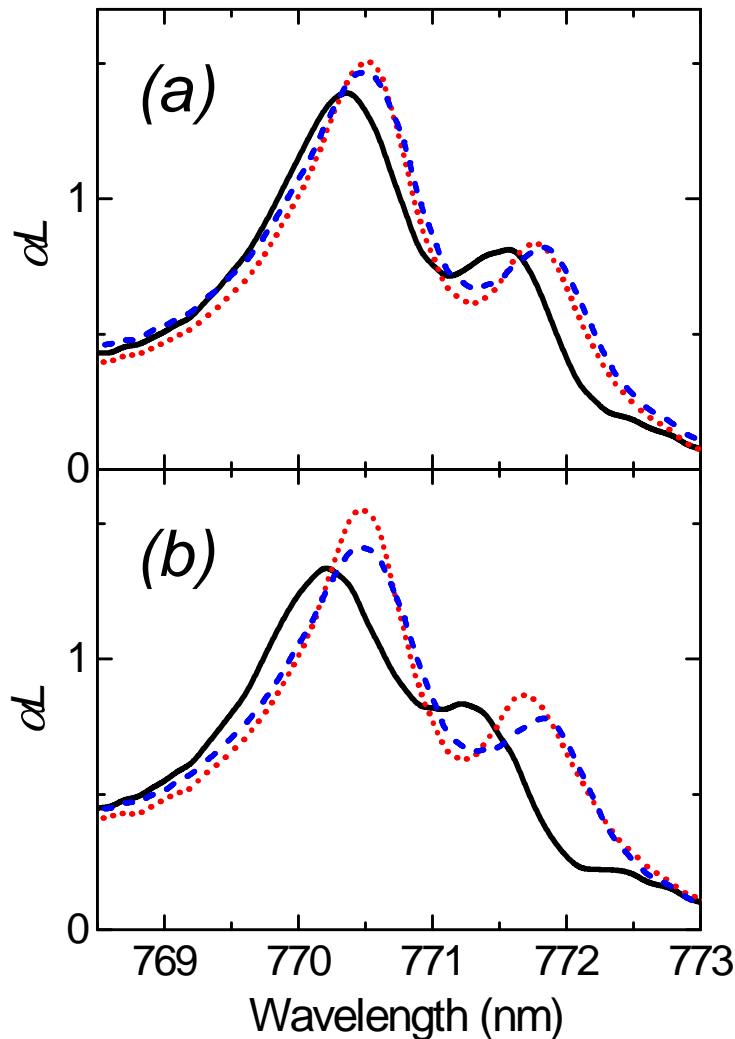
Parameters used:

$\Delta = 2$ nm, $\tau = 2$ ps (5 nm bandwidth)

Stark Shift



$I_c = 10 \text{ mW}$
 $\frac{\pi}{2}$ rotation

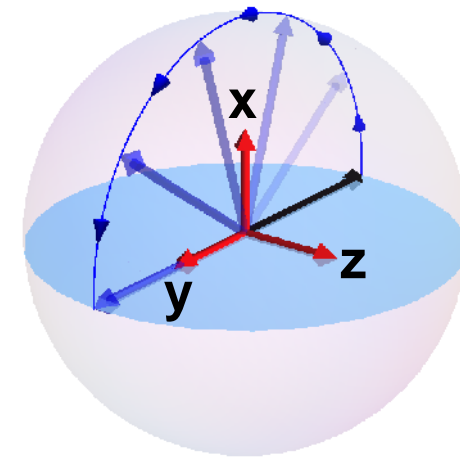
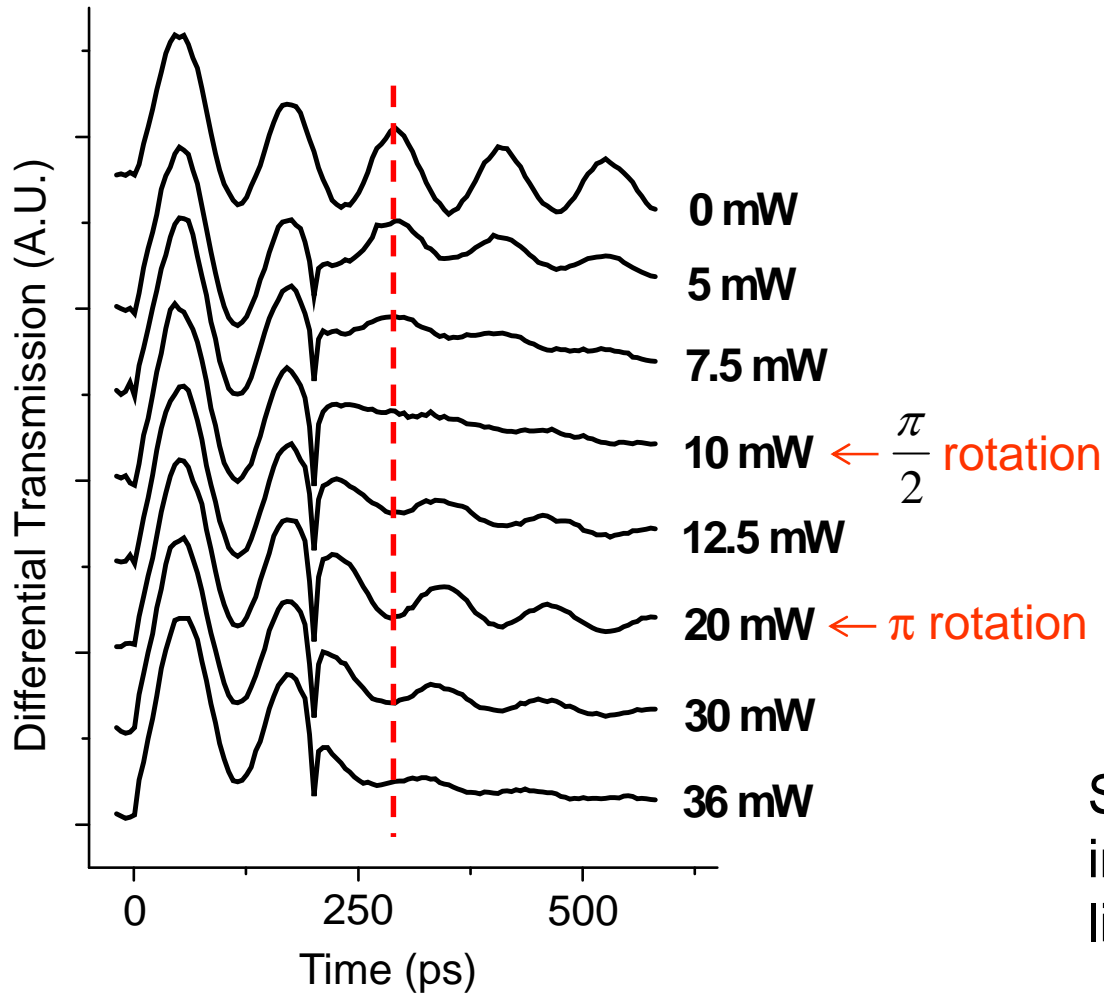
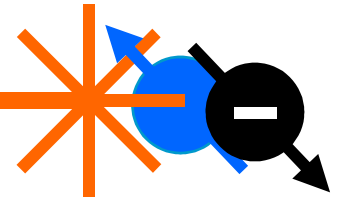


$$\phi = \int \delta\omega(t) dt$$

(a) System nearly completely recovers, leading to high fidelity rotation.

(b) Some bleaching prevents high fidelity rotation at high intensity.

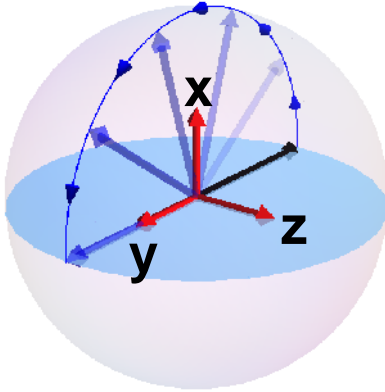
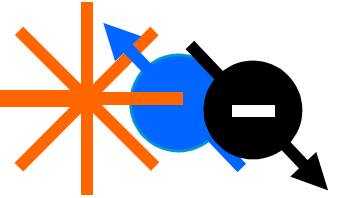
Rotation vs. Intensity



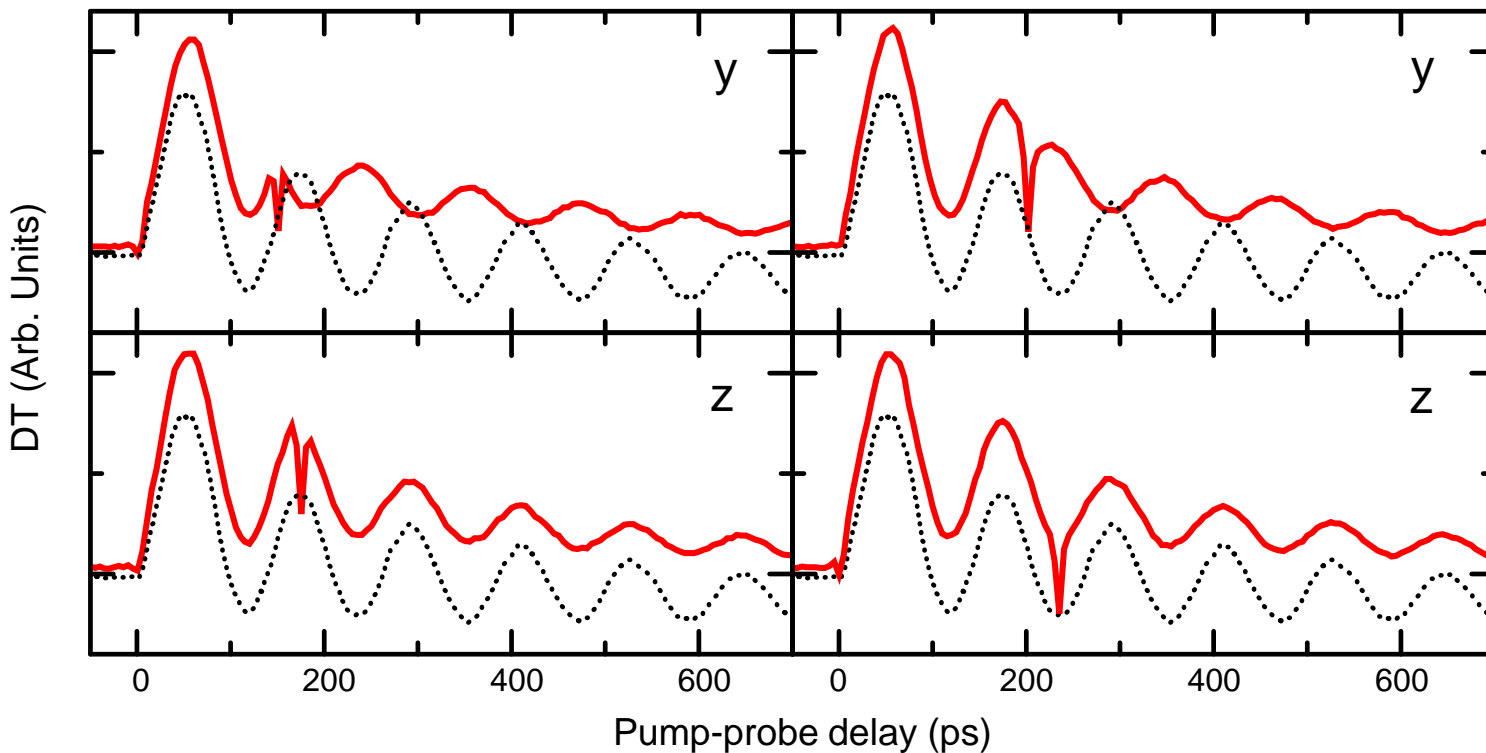
$$\delta\omega(t) \approx \frac{\Omega^2}{2\Delta}$$

Stark shift is linear for low intensity, but becomes non-linear at higher intensities.

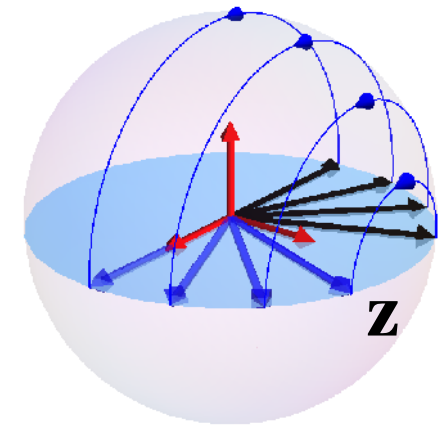
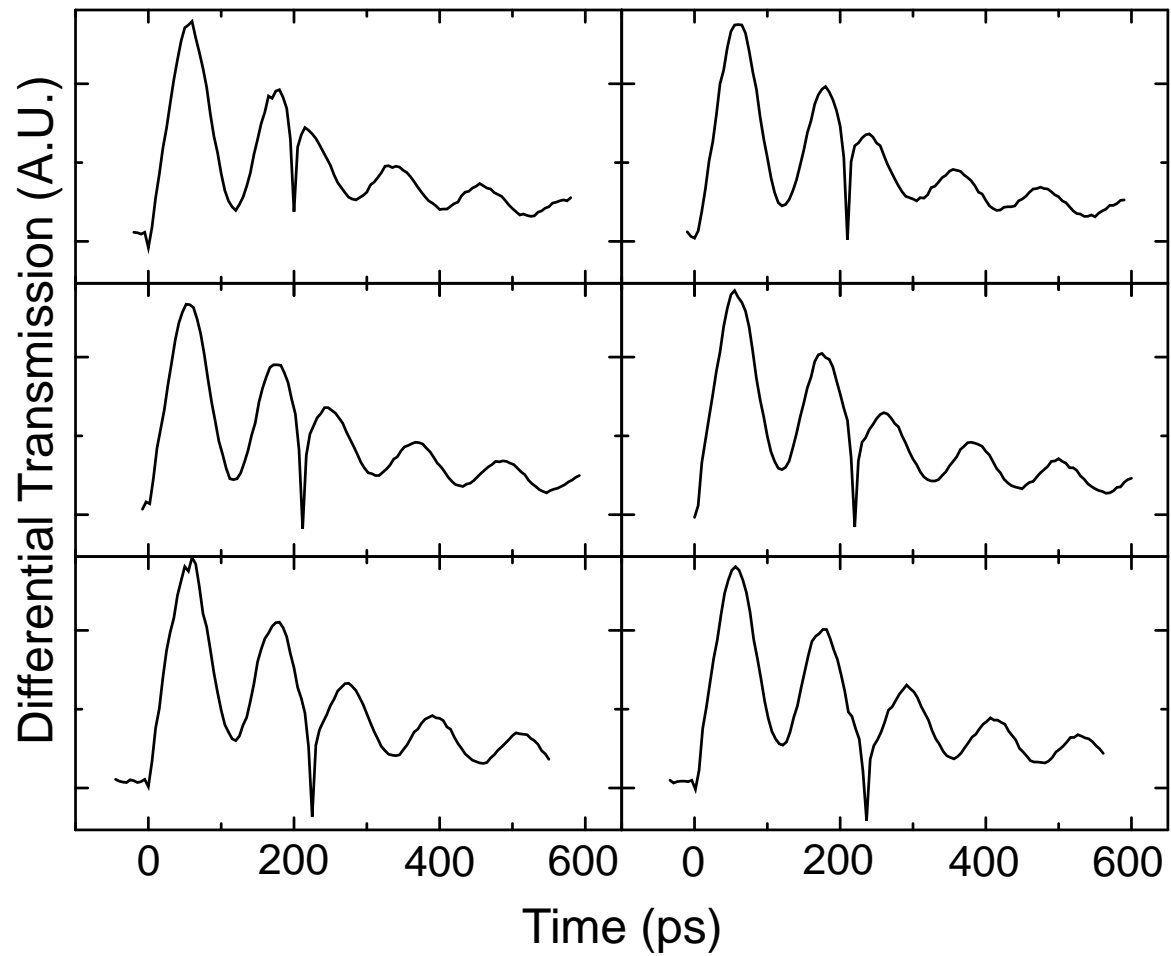
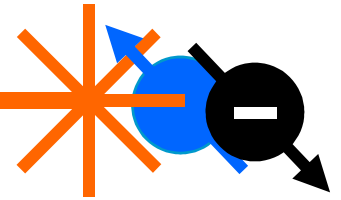
Rotation Timing



- Spins along the y-axis flip 180 degrees.
- Spins along z-axis do not rotate.

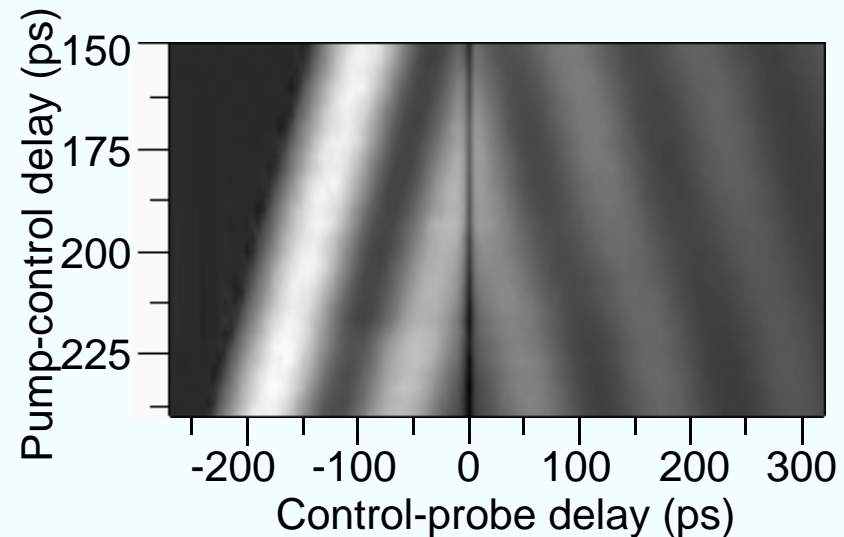
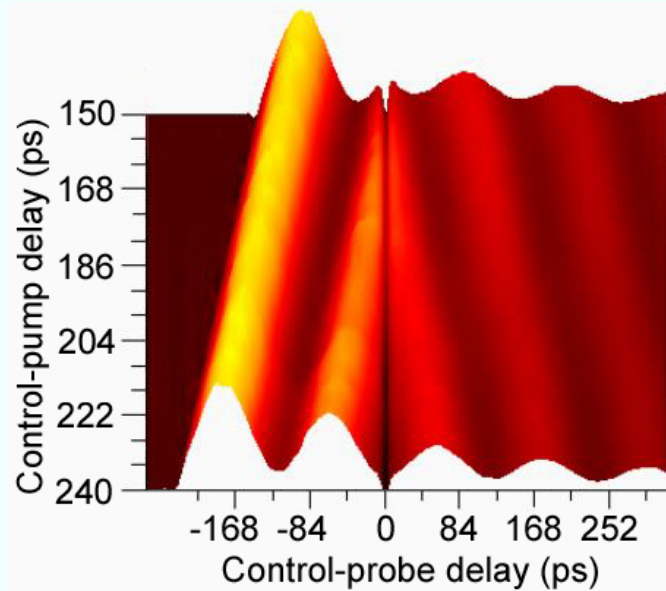
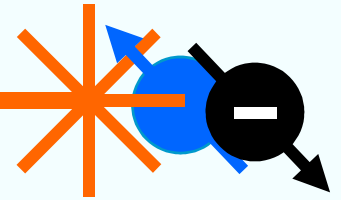


Arbitrary Timing



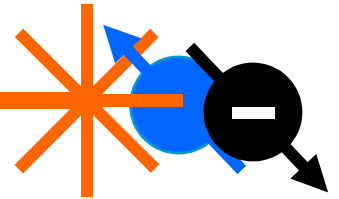
Beat peaks and troughs are symmetric w.r.t. control pulse.

Symmetrical Spin Precession



- For π rotations, phase of spin beats is symmetric w.r.t. control arrival, demonstrating complete spin flip.
- Dipole decoherence can lead to rotation asymmetry if pulse duration is too long.

Summary



- Demonstrated coherent spin flip in a 2D electron gas with an off-resonant ultrafast laser pulse, designed to induce optical Stark shifts and avoid excessive exciton excitation.
- Spin precessions are symmetric with respect to the timing of the control pulse, demonstrating the complete spin flip.
- Future work: further improvement in fidelity and decoupling of electron spins from their surroundings.